## **Chapter 4 Spatiotemporal Pattern Formation in Neural Fields with Linear Adaptation**

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**Abstract** We study spatiotemporal patterns of activity that emerge in neural fields

waves [27, 53], suggesting that some process other than inhibition must curtail

system [26, 54]. This single stationary "bump" can be perturbed and pinned with external stimuli as we see in subsequent sections of this chapter.

## **4.2.1.2 Imaginary Eigenvalues**

**Fig. 4.2** Three different cases of critical wavenumbers in the square lattice. The critical wavenumbers are (from out to in),<br> $\{\pm \atop \pm \infty\}$  $\{\pm \quad \pm \quad \},\$  $\{\pm \, , \, \pm \, - \, , \pm \, \pm \, - \, \}$  $\begin{array}{l} \{ \pm \\ \text{and} \\ \{ \pm \\ \pm \end{array} \}$  $\{\, \pm \,\, , \,\, \pm \,\, , \,\, \pm , \,\, \pm \,\, , \,\, \pm , \,\, \pm \,\, , \,\, \}$  $\pm$   $\pm$  }



see only stable traveling waves. Figure 4.1

nonzero solutions,  $z = z = z = z = 0$  which are stable if  $\leftarrow z \rightarrow z$  $\epsilon \rightarrow \epsilon + \epsilon$  -  $\rightarrow \infty$  remark that the triplet solutions  $z = z = z = 0$  *m* are never stable and that if  $\theta$  =

cycles along the principle directions. In the simulations illustrated in the figure, we change  $u_i$ .

## **4.3 Response to Inputs in the Ring Network**

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**Fig. 4.4** (a) Partition of  $(1, \alpha)$  ) parameter space into different dynamical behaviors of the bump solution (4.12) for Heaviside firing rate (4.8). Numerical simulation of the (

a moving input is introduced, the system tends to lock to it if it has speed commensurate with that of the natural wave. Converting to a wave coordinate frame  $\ell = \bullet - \bullet$  s where we choose the stimulus speed, we can study traveling wave solutions  $u \bullet s$   $\bullet \bullet s$  =  $1 \times \bullet \bullet s$  of (4.1)



**Fig. 4.5** Sloshing instability of stimulus-locked traveling bumps (4.33) in adaptive neural field (4.1) with Heaviside firing rate (4.8). (**a**) Dependence of stimulus locked pulse width  $\Delta$  on stimulus speed, calculated using the implicit equations (4.36) and (4.37). (**a**) Zeros of the Evans function E = det A<sub>p</sub>  $-1$ , with (4.47), occur at the crossings of the zero contours of Re<sup>E</sup> (*black*) and Im<sup>E</sup> (*grey*). Presented here for stimulus speed = , just beyond the Hopf bifurcation and ImE (*grey*). Presented here for stimulus speed  $=$ , just beyond the Hopf bifurcation at  $\approx$  . Breathing instability occurs in numerical simulations for (**b**)  $=$ , and (**c**)  $\approx$  . Breathing instability occurs in numerical simulations for (**b**)  $=$  **c** and (**c**)  $=$   $\cdot$   $\cdot$  (**d**) When stimulus speed  $=$   $\cdot$  is sufficiently fast, stable traveling bumps lock. Other parameters are  $\alpha = 1$ ,  $\alpha = 1$ ,  $\beta = 1$ , and  $\beta = 1$ 

 $A_{\Delta}$  =

function (4.8). In parameter regime we show, there are two pulses for each parameter value, either both are unstable or one is stable. As the speed of stimuli is decreased, a stable traveling bump undergoes a Hopf bifurcation. For sufficiently fast stimuli, a stable traveling bump can lock to the stimulus, as shown in Fig. 4.5d. However, for



dynamics of the adaptation variable  $\ell$  additionally governs the stability of the stationary bump [22]. In particular, if  $\alpha$   $\beta$ , stationary bumps are always unstable. Stable bumps in the scalar model of Amari can extend to this model only for  $\alpha$   $\beta$ , and a stable bump for  $\alpha$   $\beta$ destabilizes as  $\alpha$  decreases through  $\alpha = \beta$  leading to a drift instability [22] that gives rise to traveling bumps.

**CASE II:** *Localized Excitatory Input* / **4** . A variety of bifurcation scenarios can occur [22, 23], and, importantly, stationary bumps can emerge in a saddle-node bifurcation for strong inputs in parameter regimes where stationary bumps do not exist for weak or zero input as shown in Fig. 4.6



**Fig. 4.7** Destabilization of spatial modes  $\Omega_c$ , and  $\Omega$ , as the bifurcation parameter  $I_1$ is varied through a Hopf bifurcation, can give rise to a stable *breather* or *slosher*, respectively, depending on the relative position of the bifurcation points for each spatial mode (e.g., · and

, Fig. 4.6c). (**a**) a plot of *u*.x; t / exhibiting a breather arising from destabilization of the sum mode  $\Omega_{\text{C}}$  or  $\int$  for  $\left| \mu \right| = \sqrt{\hat{w}}$ ,  $\theta = \sqrt{\alpha}$  and  $\theta = \sqrt{\alpha}$ . (**b**) a plot of  $u \neq x$  exhibiting a slosher arising from destabilization of the difference mode  $\Omega$  or  $I_1 = \Lambda$   $\bar{w}_1 = \Lambda$  $\beta = \sqrt{\alpha} = \sqrt{1 - \frac{2}{\alpha}}$ . Common parameters:  $= \sqrt{\bar{w}} = \sqrt{1 - \frac{2}{\alpha}}$ 

the two threshold crossings of the bump relative to the position of the input  $\ell \rightarrow$ . This results in consistency conditions for the existence of a stimulus-locked traveling bump:

$$
V = \mu - \mathbf{M}_{\text{c}} \times \mu - \mu - \mathbf{M} \times \mu
$$

$$
V = \mu - \mathbf{M}_{\text{c}} \times \mu - \mu - \mathbf{M} \times \mu
$$



**Stability of Traveling Bumps.** By setting  $u = \ell$ ,  $\ell + \tilde{\ell}$  and  $\ell = \ell$ ,  $\ell + \tilde{\ell}$ , we study the evolution of small perturbations  $\tilde{C}$  <sup>T</sup> in the linearization of (4.1) about the

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