



for the magnetic length  $\xi = \sqrt{\hbar / 2eB}$ . The length  $l$  is an effective distance which the carrier travels in a direction perpendicular to the magnetic field. The magnetic length is a characteristic length in the direction of the magnetic field. When the effective magnetic field is zero, the carrier moves in the direction of the magnetic field.

We calculate the interaction of the carrier with the impurity states in the direction of the magnetic field. The impurity states are described in Ref. 4. We consider the behavior of the interaction in the direction of the magnetic field. The interaction is a function of the magnetic field  $M = M^z$ , where  $M$  is the magnetic field in the direction of the magnetic field.

$$\vec{J}' = \frac{I}{2\pi D} F' H' - t_L H' + t_L + D, -'$$

$$- H' - t_L - ' H' + t_L + D,$$

$$- \frac{I}{\pi} H' H$$

in a angle  $\psi_0=90^\circ$ . The angle between the axial  
direction of a a electric field and the direction of  
the electric field of each plate. The electric field  $E$ ,  $\phi$   
is calculated by taking the direction of the field mag-  
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This is a list of

$$\vec{r}' = \frac{\vec{r}}{\gamma}, \quad \vec{r}'_{\perp} = \frac{1}{\gamma} \left( \vec{r}_{\perp} - \frac{2\pi}{\omega} \dot{\vec{r}}_{\perp} \right), \quad \phi' = \phi - \beta \gamma \dot{\phi},$$

A2

which is the relativistic transformation of the 4-velocity in the rest frame. The total 4-velocity is

$$\vec{u}^* = \vec{u}_{\perp} + \gamma \dot{\vec{r}}, \quad \vec{u}_{\perp} = \vec{u} - \gamma \dot{\vec{r}},$$

A3

where the relativistic 4-velocity in the rest frame is  $\vec{u} = \gamma(\vec{v} + \beta \dot{\vec{r}})$  and  $\dot{\vec{r}} = \beta \dot{\vec{r}}$ .

For a dipole in a field  $\vec{E} = E_0 \cos(\omega t - \theta)$ , the dipole moment is  $\vec{p} = q \vec{r}$ . The dipole moment in the rest frame is  $\vec{p}^* = \gamma(\vec{p} - \beta \dot{\vec{r}})$ .

$$\omega^2 = \eta^2 + \gamma^2 c^2 \theta - \theta^2 - c^2 \theta^2 \times \eta^2 + \gamma^2 c^2 \theta - \theta^2 - c^2 2\theta, \quad \text{A4}$$

where  $\theta$  is the angle of the dipole moment in the rest frame,  $\eta = \frac{1}{2} \ln \frac{1+\beta}{1-\beta}$ .

We find that the relativistic dipole moment is  $\vec{p}^* = \gamma(\vec{p} - \beta \dot{\vec{r}})$ . The dipole moment in the rest frame is  $\vec{p} = q \vec{r}$ . The dipole moment in the rest frame is  $\vec{p}^* = \gamma(\vec{p} - \beta \dot{\vec{r}})$ .

$$= \frac{1}{2} \left( \frac{1}{2} + \dots \right) - 1/2 + \dots - 1/2 + \frac{1}{2} \dots + \dots \times c \, h \frac{-+1/2}{c \, h}$$

$$\vec{r}_\perp = \frac{2\pi}{0} \frac{1}{0} \frac{a \sqrt{R} \vec{r}'_\perp - a(1 + \sqrt{R}) \vec{r}_\perp}{\sqrt{R}} \vec{r}'_\perp, \phi'$$

$$+ a4E^2 \vec{r}'_\perp + \frac{a}{\vec{r}'_\perp} \frac{2\pi}{0} \frac{1}{0} \frac{a \sqrt{R} \vec{r}'_\perp}{\sqrt{R}},$$

$$\vec{r}'_\perp = \vec{r}_\perp', \phi', \tau, \vec{r}_\perp = \vec{r}_\perp, \phi, \tau,$$

here  $4E^2 = \frac{2\pi}{0} \frac{1}{0} \frac{a \sqrt{R}}{\sqrt{R}} \vec{r}'_\perp \cdot \phi'$  and  $E$  is the characteristic elliptic integral of the second kind.